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# In GaAs/GaAs/In GaP Strained Quantum Well Lasers Grown by Metalorganic Chemical Vapor Deposition

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**Abstract** A lum inum-free strained Ino  ${}_{2}$ Gao  ${}_{8}$ A s/GaA s/Ino  ${}_{4}$ 9Gao  ${}_{51}$ P single quantum well lasers with an emission wavelength of 985nm are presented. The laser material was grown by metalorganic chemical vapor deposition. An extremely low threshold current density of 150A/cm² is obtained for  $100\mu$ m wide stripe lasers having a cavity length of  $800\mu$ m. The internal quantum efficiency and the internal loss are 78% and 5cm², respectively. The measured vertical beam divergence angle from the laser is about 45.°

PACC: 4255P, 6855, 7865, 8115H

## 1 Introduction

Strained InGaA s/GaA s quantum well (QW) lasers with emission wavelength of 980nm are currently receiving considerable attention because they yield a lower noise figure, higher gain coefficient than 1. 48µm InGaA sP laser in the application of the erbiumdoped fiber amplifier (EDFA) pumping sources<sup>[1,2]</sup>. In addition, the InGaA s/GaA s strained quantum well lasers have lower threshold current and high slope efficiency. A l-GaA s is usually used as cladding layers in most of the InGaA s/GaA s QW lasers But in recent years, a luminum-free materials for the claddings are used to improve the reliability of InGaA s/GaA s strained quantum well lasers. In 49Ga 51P lattice matched with GaA s, was introduced as a substitution for A IGaA s cladding layers grown by low-pressure metalorganic chemical vapor deposition (LPMOCVD) [3~6] or gas-source molecular beam epitaxy (GSMBE) [7~10].

The advantages of InGaA s/GaA s/InGaP QW lasers are as follows Since it is aluminum-free, less surface oxidation during the fabrication process and laser operation is ex-

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pected The low surface recombination velocity will enhance catastrophic optical damage (COD) threshold of the laser facets and improve the reliability of the lasers. Further more, the successfull selective chemical etching between GaAs and InGaP layers makes the ridge-waveguide structure processing much more easily controlled. In this paper, we report on the strained InGaAs/GaAs/InGaP single quantum well lasers grown by MOCVD.

## 2 Laser Structure and Material growth

The substrate used was Si-doped (100)  $n^+$ -GaAs wafers, tilted 6 °towards the 111 A plane, with the etch pit densities of less than 500cm<sup>-2</sup>. The layer structure of the In-GaAs/GaAs/InGaP laser is as follows (see Fig. 1): a  $0.3\mu$ m-thick  $n^+$ -GaAs buffer layer, a  $1.0\mu$ m-thick  $n^-$ In<sub>0</sub> 49Ga<sub>0</sub> 51P lower cladding layer, an undoped  $0.1\mu$ m-thick GaAs lower waveguide layer, an undoped  $0.1\mu$ m-thick GaAs as the upper waveguide layer, a  $1.0\mu$ m-thick p-In<sub>0</sub> 49Ga<sub>0</sub> 51P upper cladding layer, a  $0.1\mu$ m-thick p<sup>+</sup>-GaAs cap layer, and a 30nm heavily doped p<sup>++</sup>-GaAs ohm ic contact layer for minimizing the contact resistance

The laser material was grown by low-pressure MOCVD in an A IXTRON-200 system at a constant temperature of 700, except for the top GaAs layers grown at 600 to increase the Zn incorporation. The source materials were trimethylgallium (TMG), trimethylindium (TMI), arsine (AsH3) and phosphine (PH3). Dimethylzinc (DMZ) and silane (Si2H8) were used as the n- and p-type dopants, respectively.

Since the layer structures contain relatively thick InGaP cladding layers, it is essential to have close lattice matching to

$p^{++}$ -GaA s ohm ic contact (Zn: 1 × 10 <sup>20</sup> cm <sup>-3</sup> )
$p^+$ -GaA s cap (Zn: 1 × 10 <sup>19</sup> cm - <sup>3</sup> )
p-InGaP upper clad (Zn: $1 \times 10^{18}$ cm <sup>-3</sup> )
i GaA s upper w aveguide (undoped)
i- InGaA s quantum well (undoped)
i- GaA s lower waveguide (undoped)
n-InGaP lower clad (Si: $1 \times 10^{18}$ cm <sup>-3</sup> )
$n^+$ -GaA s buffer (Si $1 \times 10^{18}$ cm $^{-3}$ )
n <sup>+</sup> -GaA s substrate (Si: $2 \times 10^{18}$ cm <sup>-3</sup> )

Fig. 1 The layer structure of the InGaA s/GaA s/InGaP laser

the GaAs substrate It is experimentally observed that the lattice mismatch  $(\Delta a/a)$  of the relatively thick  $(1.0\mu\text{m})$  In GaP layers must be less than  $2 \times 10^{-3}$  to avoid degradation in laser performance

### 3 Device Fabrication and Laser Performance

The laser structure is evaluated by fabrication and characterization of wide stripe  $(100\mu\text{m} \text{ wide})$  devices. The devices are fabricated by chemical etching through the p<sup>+</sup> - and p<sup>++</sup>-GaAs layers outside the  $100\mu\text{m}$  stripe to prevent current spreading. A 120nm-thick  $SO_2$  was deposited over the entire p-surface by Plasm a-Enhanced Chemical V apor Deposition.

tion (PECVD). Then the wafer was processed into  $100\mu m$  wide broad stripes by using standard photolithography techniques. After the p-contact consisting of Ti/Pt/Au was deposited, the wafer was thinned to a thickness of about  $100\mu m$ , then the n-contact consisting of AuGe/Ni/Au was deposited. The lasers were then cleaved into bars with various cavity lengths. The tested devices were mounted p-side down on a copper heat-sink using indium solder.

Figure 2 shows the plot of the threshold current density against various cavity lengths for the fabricated lasers. The threshold current density is  $150 \text{A}/\text{cm}^2$  at a cavity length of  $800 \mu \text{m}$ . When the cavity length extends to  $1200 \mu \text{m}$ , the threshold current density decreases to  $130 \text{A}/\text{cm}^2$  due to the decreased facet optical loss

From the relationship of the reciprocal of the external quantum efficiency versus the cavity length for the fabricated lasers, an internal quantum efficiency and total internal loss of 78% and 5cm<sup>-1</sup> are deduced, respectively. It is expected that the use of InGaA sP (> 1. 5eV) waveguide layer or a step graded index separate confinement heterostructure will improve the carrier injection efficiency, and consequently enhance the internal quantum efficiency compared with the GaA swaveguide layer structure shown in Fig. 1.

From the perpendicular far-field intensity profile, we know that the beam divergence or far-field angle at full width of half maximum (FW HM) is about 45.° The large vertical beam divergence is attributed to the strong optical confinement in the laser structure. To meet the demand of the device in system applications, it would be necessary to reduce the beam divergence so that high fiber coupling efficiency could be achieved. This can be obtained by reducing the optical confinement in the single quantum well laser structure. It is very important to design the waveguide layer structure with narrow far-field angle and low threshold current.

Figure 3 gives the lasing spectrum measured at 200mW for the  $100\mu m \times 800\mu m$  broad stripe laser device. The emitting wavelength is 985nm and the FW HM is 3nm.

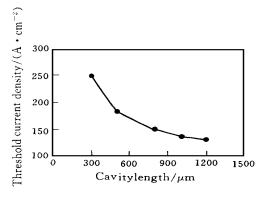


Fig 2 Threshold current density of InGaA s/GaA s/InGaP strained quantum well lasers as a function of cavity length

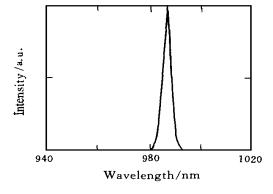


Fig 3 M easured lasing spectrum for the InGaA s/GaA s/InGaP laser device

The initial aging tests for the A l-free lasers have demonstrated better reliability than the A l-containing lasers. The detail of the aging test results and the improvement in laser performance will be reported elsewhere

## 4 Conclusion

In conclusion, we report on the room temperature CW operation of strained In<sub>0</sub> 2Ga<sub>0</sub> 8 A s/GaA s/In<sub>0</sub> 49 Ga<sub>0</sub> 51 P single quantum well lasers grown by MOCVD. The broad area lasers show extremely low threshold current density of  $150A/cm^2$  at a cavity length of  $800\mu$ m. Our results demonstrate that MOCVD is also suitable for growing the 980nm In-GaA s/GaA s lasers with In<sub>0</sub> 49Ga<sub>0</sub> 51P as the cladding layers

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